

# Probing the Neutrino Mass Hierarchy in INO using the Atmospheric $\nu_\mu$ Survival Rate

# Raj Gandhi



Harish-Chandra Research Institute, Allahabad, India

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- Conclusions

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- Precision allows us to identify or exclude special valued mixing angles like  $\theta_{13} = 0^\circ$ ,  $\theta_{23} = 45^\circ$ , and special relations between the quark and lepton sectors like  $\theta_{12} + \theta_C = 45^\circ$ , check for unitarity of 3 generations, non-standard interactions, decoherence scenarios.....

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- A multi-pronged effort to achieve these goals is underway via various operating, planned and proposed experiments, e.g. long-baseline, reactor, atmospheric, solar, beta-decay,  $\nu$ -less double beta decay, precision cosmology etc.

## *Introduction . . .*

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- $(\nu_\alpha) = U(\nu_i)$  with  $i = 1, 2, 3$  and  $\alpha = e, \mu, \tau$ .
- The  $3 \times 3$  neutrino mixing matrix  $U$  in the MNS parametrization is:

$$\begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

where  $c_{ij} = \cos \theta_{ij}$ ,  $s_{ij} = \sin \theta_{ij}$  ( $\theta_{ij}$  = mixing angle between  $i^{th}$  &  $j^{th}$  mass states).  $\delta$  = CP phase.

## *Introduction . . .*

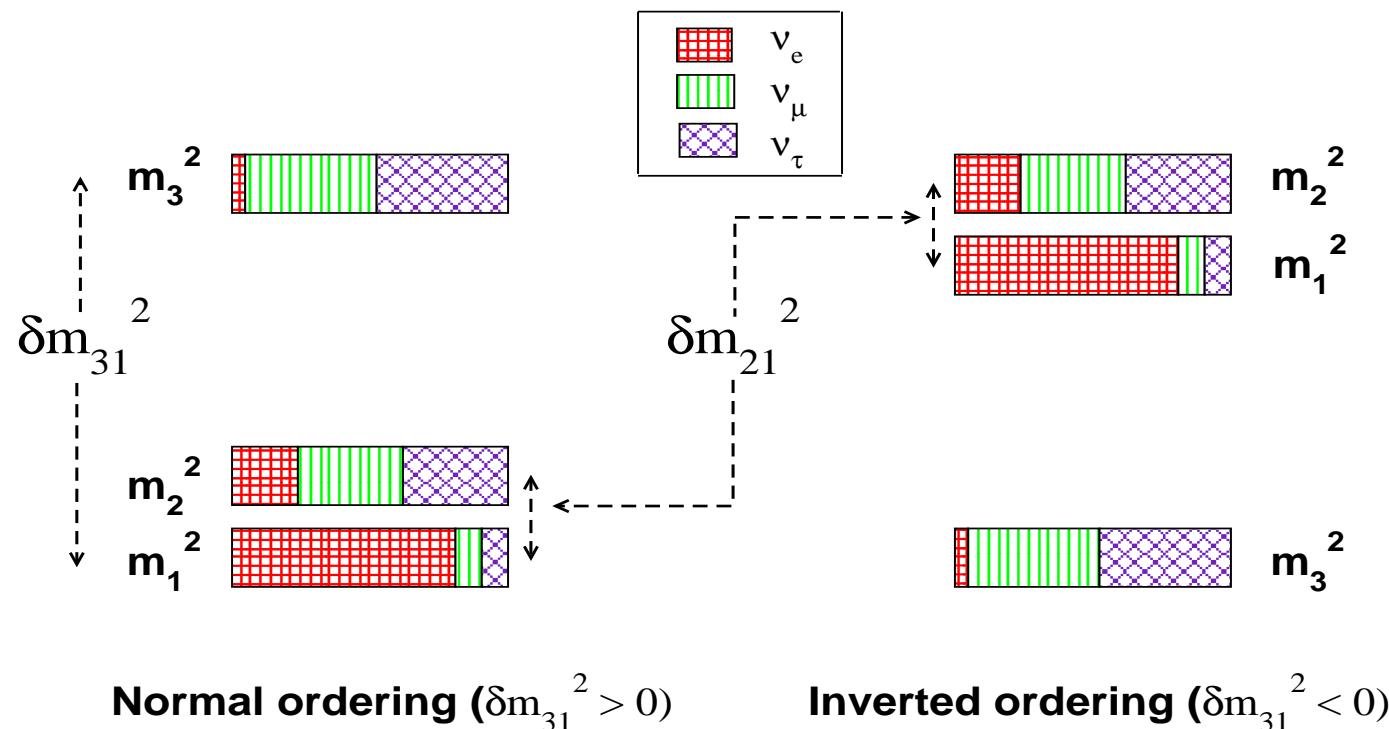
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- $\tan^2 \theta_{12} = 0.39^{+0.05}_{-0.04}$  solar,reactor
- $\tan^2 \theta_{23} = 1.0^{+0.3}_{-0.3}$  atmospheric, K2K
- $\theta_{13} < 12^\circ @ 3\sigma$  reactor, atmospheric
- $\delta m_{21}^2 = 8.2^{+0.3}_{-0.3} \times 10^{-5} \text{ eV}^2$  solar,reactor
- $|\delta m_{31}^2| = 2.2^{+0.6}_{-0.4} \times 10^{-3} \text{ eV}^2$  atmospheric, K2K
- $m_\beta < 2.2 \text{ eV}$  beta decay
- $m_{\beta\beta0\nu} < 0.3 \text{ eV}$   $\nu$  less double beta decay
- $\sum m_i < 1.6 - 0.7 \text{ eV}$  Precision Cosmology
- $\delta_{CP}$  is unknown

## $\nu$ The Mass Hierarchy, its Significance and Detection

So far we only know  $|\Delta m_{31}^2|$  and not its Sign



**Normal ordering** ( $\delta m_{31}^2 > 0$ )

**Inverted ordering** ( $\delta m_{31}^2 < 0$ )

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- A large class of GUTS use the Type I seesaw mechanism to unify quarks and leptons. Several positive features are lost if in such models the neutrino hierarchy is *inverted* rather than *normal*

C. Albright, Phys. Lett. B599, 285 (2004)

- An *inverted* hierarchy on the other hand would generally require  $m_1, m_2$  to be *degenerate*, hinting towards an additional global symmetry in the lepton sector.

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- It would also favour theories utilising *the Type II seesaw mechanism with additional Higgs triplets*.
- The type of hierarchy impacts the effectiveness of *leptogenesis* in most theoretical models.

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  - **Nu Factories** substantial improvement in bounds

- Detection of the hierarchy via the  $\nu_\mu \rightarrow \nu_e$  channel, at not-so-long baselines, is hampered by a  $\delta_{CP} - \text{sign}(\Delta_{31})$  degeneracy which reduces sensitivity. Additionally, in “off-resonance” situations, there is a  $\theta_{13} - \text{sign}(\Delta_{31})$  that must be taken into account.

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- Overcoming this in superbeam experiments would require using either *multiple oscillation channels, 2 baselines, different energies, the magic baseline, 2 off-axis locations* or a combination of these strategies.

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- These degeneracies are absent in methods which utilise the muon survival probability.

- Vacuum 1-MSD limit ( $\Delta m_{21}^2 = 0$  ) :

$$\mathcal{P}_{\mu e} = \sin^2 \theta_{23} \sin^2(2\theta_{13}) \sin^2 [\Delta_{31} L/E]$$

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- and  $\Delta_{31}^m = \Delta_{31} \sqrt{(\frac{A}{\delta m_{31}^2} - \cos 2\theta_{13})^2 + \sin^2 2\theta_{13}}$

- We note that:  $\mathcal{P}_{\mu e}^m$  is maximized not just at resonance but when the combination  $\sin^2(2\theta_{13}^m) \sin^2 [\Delta_{31}^m L/E]$  is maximal, *i.e.* when  $E = E_{res} = E_{peak}^m$ .

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- The  $L$  where these ( $p = 0$ ) maxima occur are 7600 km for  $\sin^2 2\theta_{13} = 0.2$ , 10200 km for  $\sin^2 2\theta_{13} = 0.1$  and 11200 km for  $\sin^2 2\theta_{13} = 0.05$ .

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- *Co-relations and degeneracies play an important role and reduce sensitivity when  $\mathcal{P}_{\mu e}^m$  is used to determine the hierarchy*

## Matter effect on $P(\nu_\mu \rightarrow \nu_\tau)$

- The 1-MSD vacuum expression is:

$$\mathcal{P}_{\nu_\mu \rightarrow \nu_\tau}^v = \cos^2 \theta_{13} \sin^2 2\theta_{23} \sin^2(\Delta_{31})$$
$$- \cos^2 \theta_{23} P_{\nu_\mu \rightarrow \nu_e}^{vac}$$

- The 1-MSD analytic expression in matter is:

$$\mathcal{P}_{\nu_\mu \rightarrow \nu_\tau}^m = \cos^2 \theta_{13}^m \sin^2 2\theta_{23} \sin^2[(\Delta_{31} + A + \Delta_{31}^m)/2]$$
$$+ \sin^2 \theta_{13}^m \sin^2 2\theta_{23} \sin^2[(\Delta_{31} + A - \Delta_{31}^m)/2]$$
$$- \cos^2 \theta_{23} P_{\nu_\mu \rightarrow \nu_e}^{mat}$$



## Maximizing the matter effect in $\mathcal{P}_{\mu\tau}$

- $E = E_{res} = E_{peak}^v$  leads to:

$$\Delta\mathcal{P}_{\mu\tau} \simeq \cos^4[\sin 2\theta_{13}(2p+1)\pi/4] - 1$$

(with  $\sin^2 2\theta_{23} = 1$ ,  $\cos^2 \theta_{13}, \cos 2\theta_{13} \simeq 1$ ).

- The maximum matter effect condition for  $L$  is:

$$[\rho L]_{\mu\tau}^{max} = (2p+1)\pi \times 5.18 \times 10^3 \times \cos 2\theta_{13} \text{ km gm/cc}$$

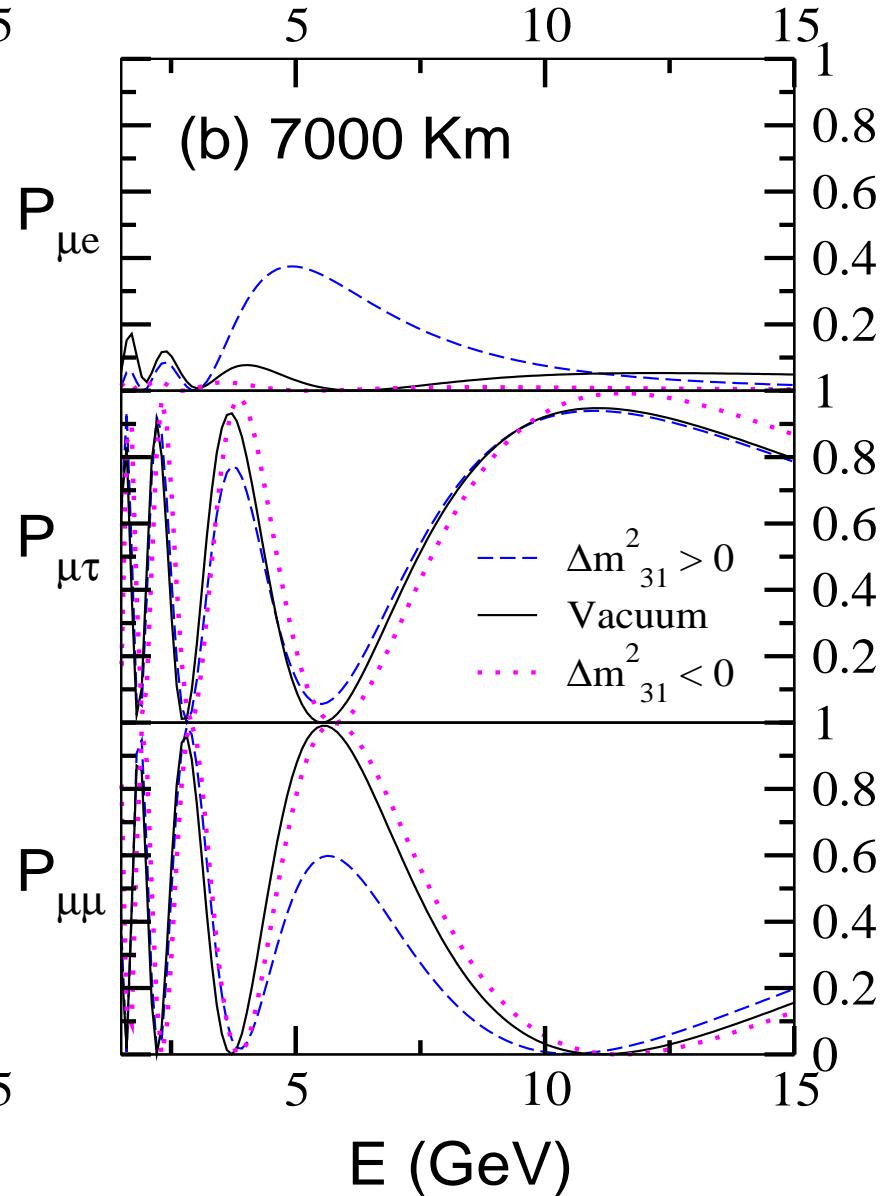
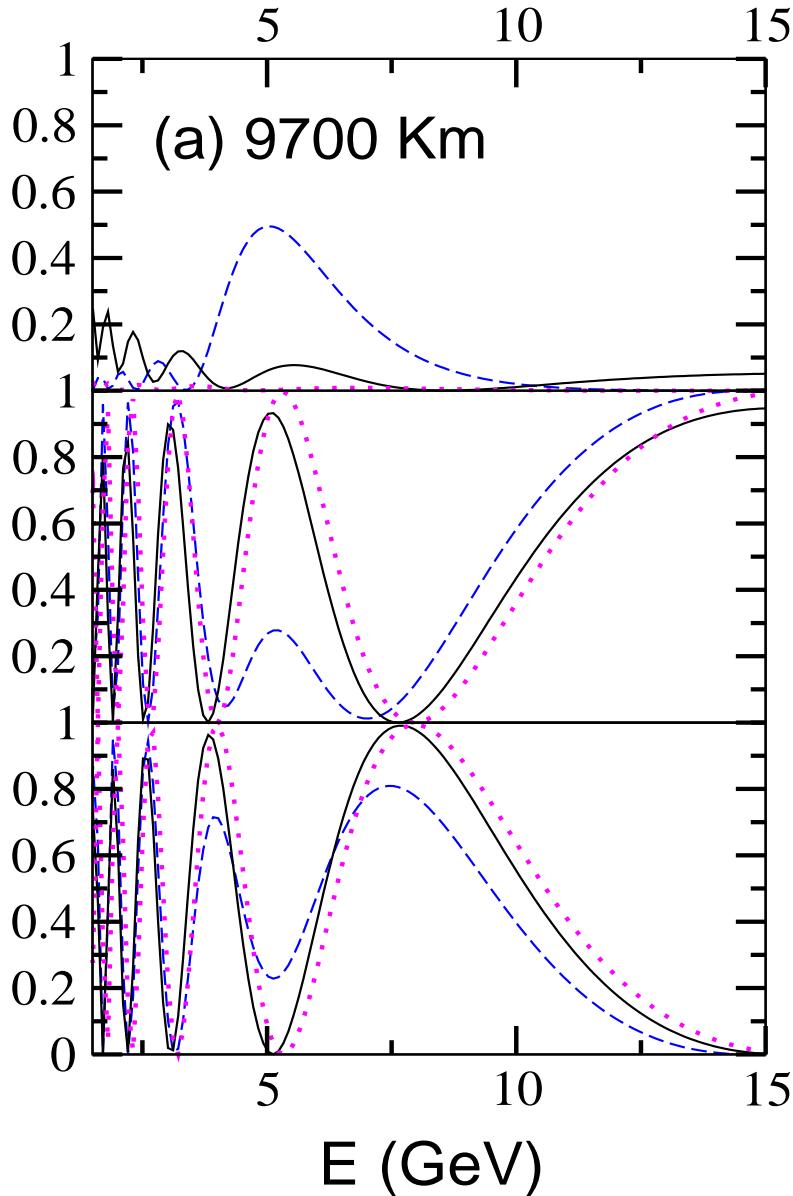
The maximum matter effect occurs for  $L = 9700$  km for  $p=1$  &  $\sin^2 2\theta_{13} = 0.1$  (9300 km, 9900 km for 0.2, 0.05).

Thus, conditions for maximizing the matter effects at very long baselines are:

- $[\rho L]_{\mu e}^{\max} = \frac{(2p+1)\pi 5.18 \times 10^3}{\tan 2\theta_{13}}$  Km gm/cc.
- $[\rho L]_{\mu\tau}^{\max} = (2p + 1)\pi 5.18 \times 10^3 \cos 2\theta_{13}$  Km gm/cc.
- $[\rho L]_{\mu\mu}^{\max} = p\pi \times 10^4 \cos 2\theta_{13}$  Km gm/cc.

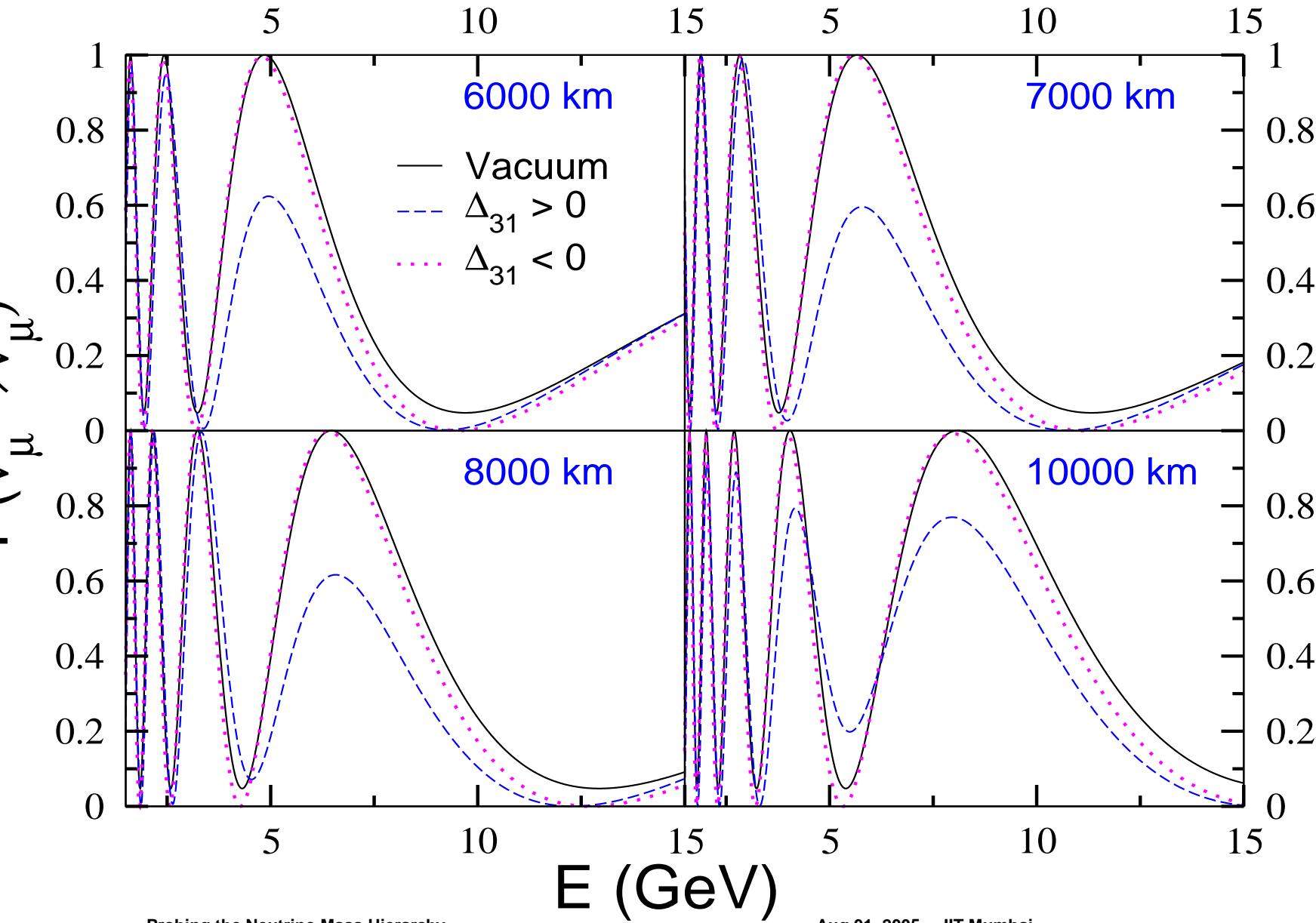
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R. Gandhi *et al.*, Phys. Rev. Lett. 94, 051801 (2005); hep-ph/0411252

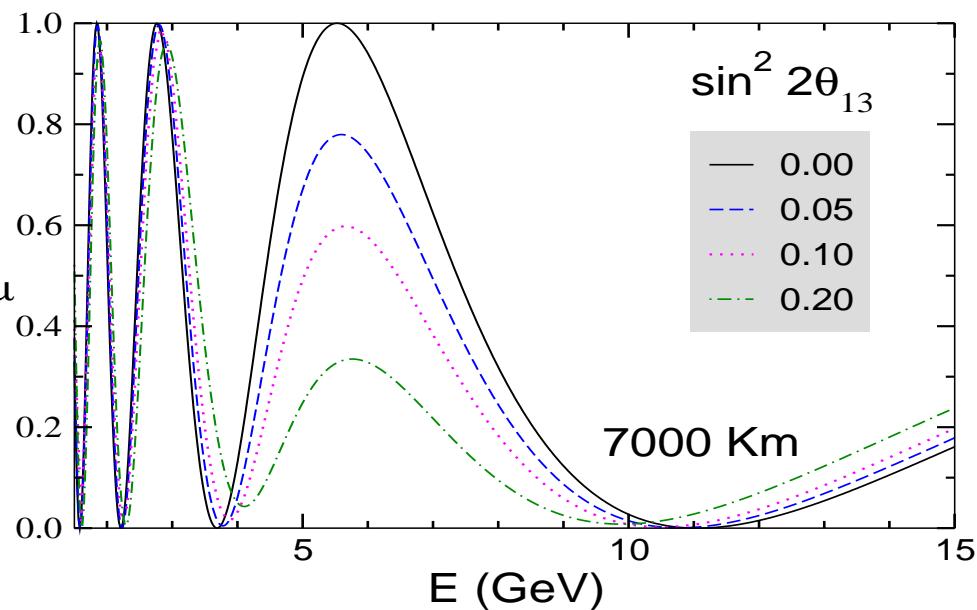


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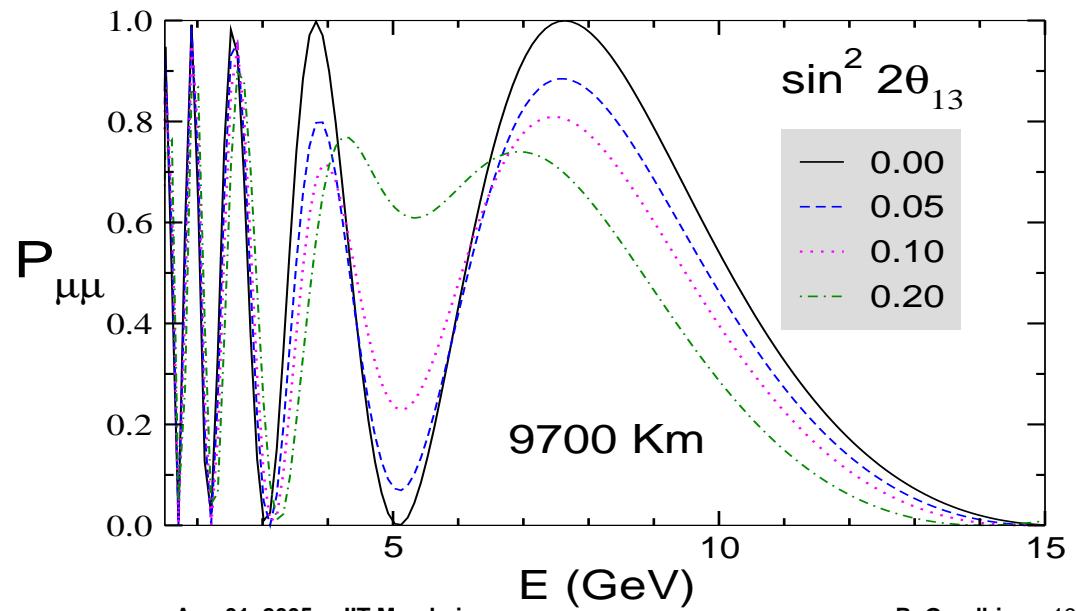
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# Survival Probability : $\theta_{13}$ sensitivity . . .



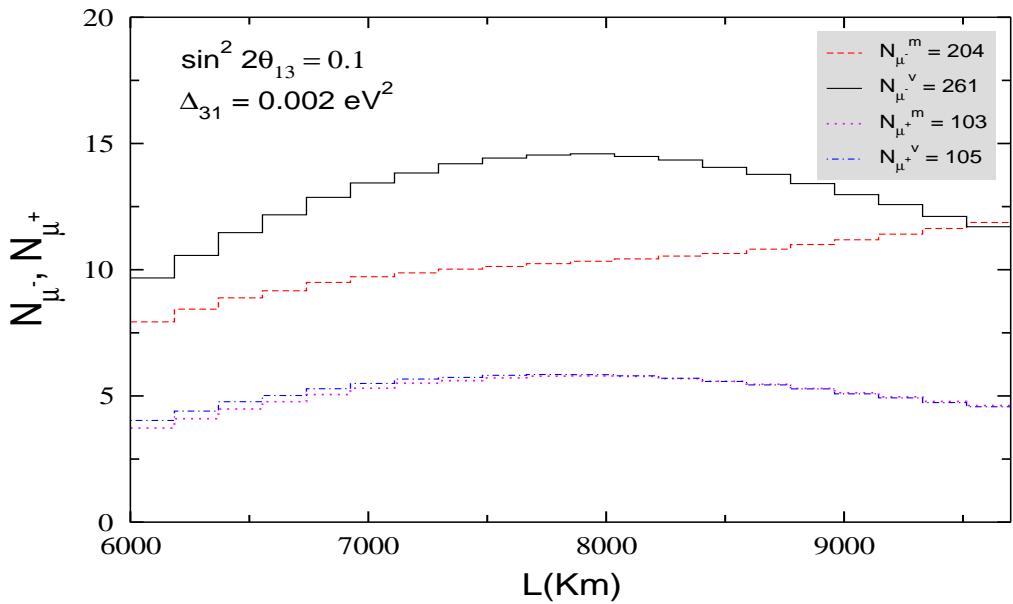
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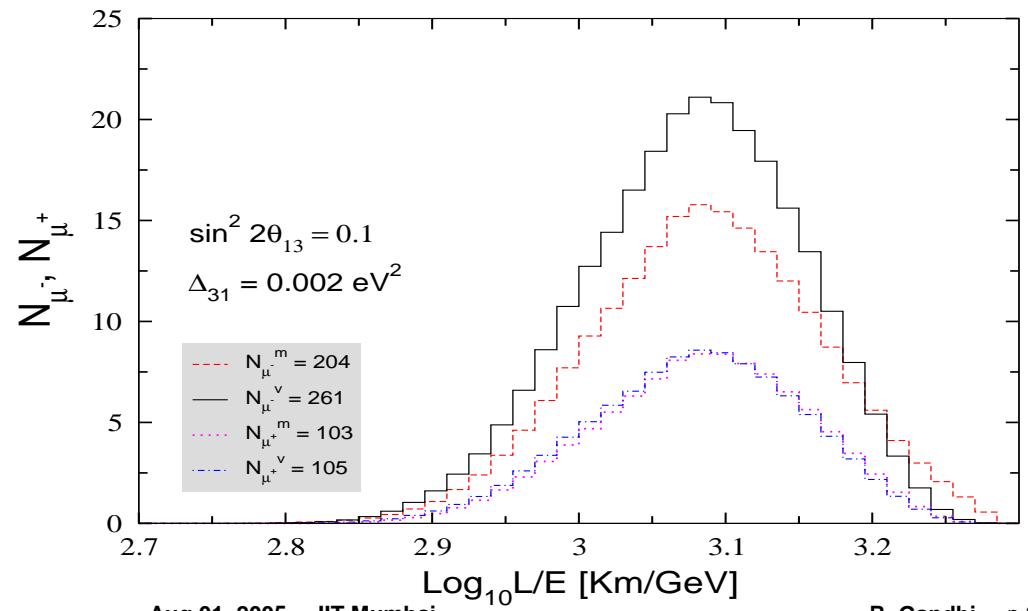
## *Results for an Iron Calorimeter type detector*

# Results : Iron Calorimeter, 1000 kt-yr . . .

$L = 6000$  to  $9700$  Km,  $E = 5$  to  $10$  GeV



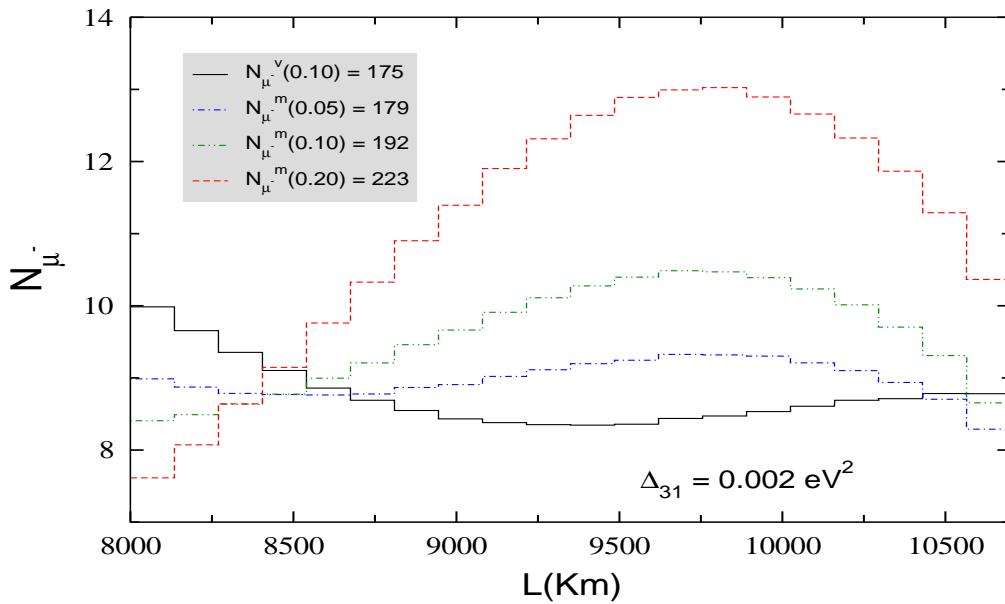
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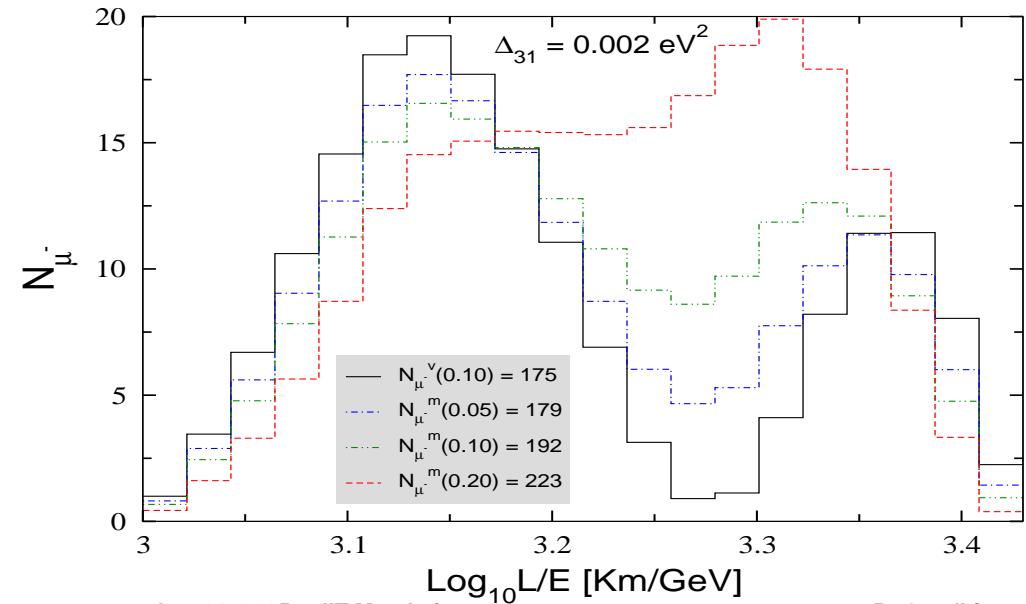
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$L = 8000 \text{ to } 10700 \text{ Km, } E = 4 \text{ to } 8 \text{ GeV}$



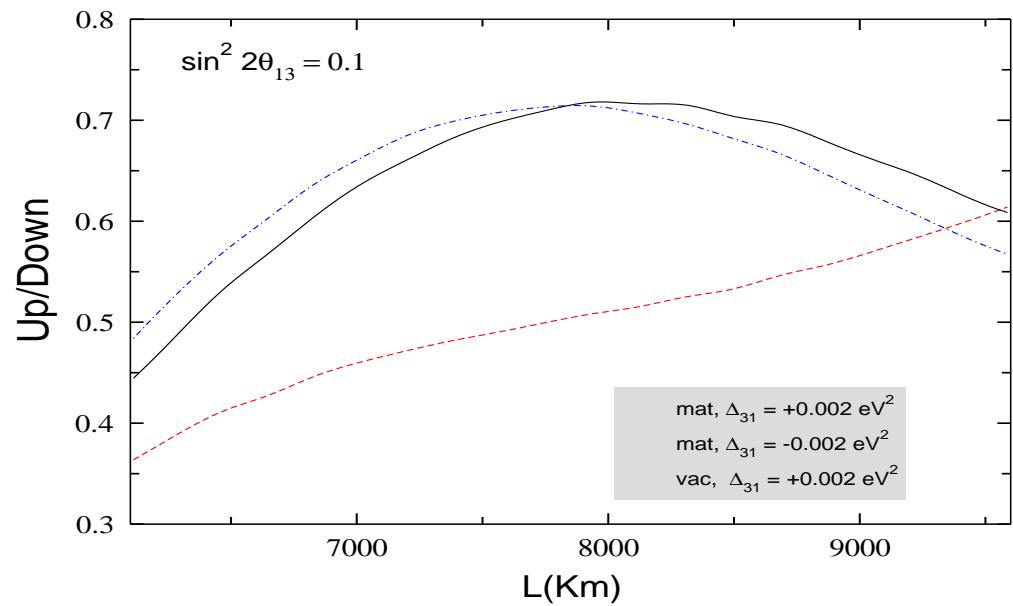
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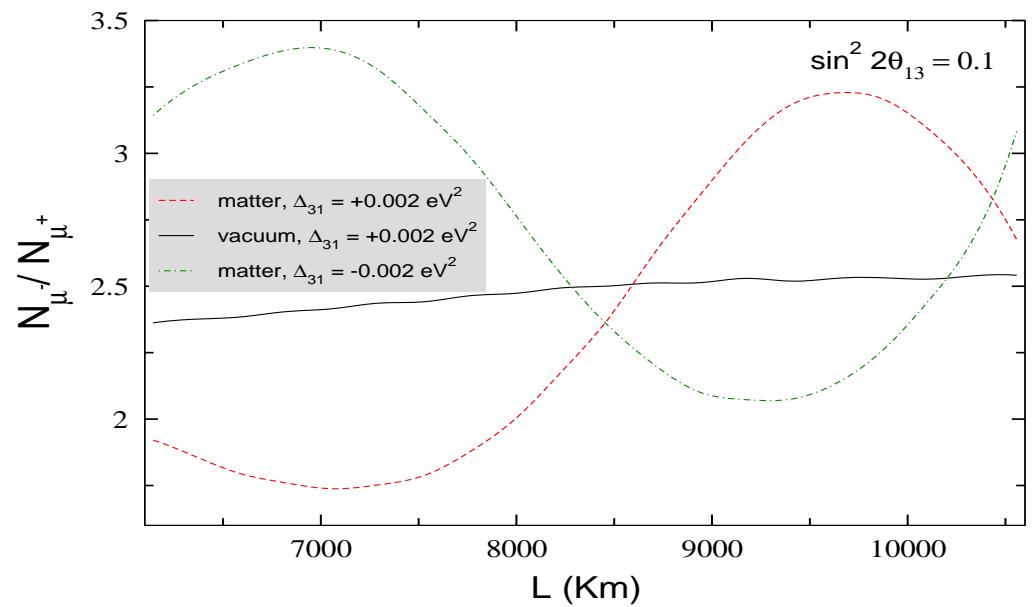
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Choice of the optimal ranges for  $\mu^-$  events for 1000 kiloton-yr:

Range 1 :  $E = 5 - 10$  GeV and  $L = 6000 - 9700$  Km

| $\sin^2 2\theta_{13}$ | $N_{vac}$ | $N_{mat}^{NH}$ | $N_{mat}^{IH}$ | $\sigma_{NH-vac}$ | $\sigma_{NH-IH}$ |
|-----------------------|-----------|----------------|----------------|-------------------|------------------|
| 0.05                  | 260       | 227            | 264            | $2.2\sigma$       | $2.5\sigma$      |
| 0.1                   | 261       | 204            | 262            | $4.0\sigma$       | $4.1\sigma$      |
| 0.2                   | 263       | 163            | 261            | $7.8\sigma$       | $7.7\sigma$      |

| $\sin^2 2\theta_{13}$ | $N_{vac}$ | $N_{mat}^{NH}$ | $N_{down}$ | $\sigma_{NH-vac}$ |
|-----------------------|-----------|----------------|------------|-------------------|
| 0.05                  | 260       | 227            | 410        | $1.8\sigma$       |
| 0.1                   | 261       | 204            | 410        | $3.3\sigma$       |
| 0.2                   | 263       | 163            | 410        | $6.6\sigma$       |

Range 2 :  $E = 4 - 8$  GeV and  $L = 8000 - 10700$  Km  
 $(\log_{10}(L/E) = 3.21 - 3.44)$

| $\sin^2 2\theta_{13}$ | $N_{vac}$ | $N_{mat}^{NH}$ | $\sigma^{NH-vac}$ |
|-----------------------|-----------|----------------|-------------------|
| 0.05                  | 23        | 43             | $3.0\sigma$       |
| 0.1                   | 23        | 63             | $5.0\sigma$       |
| 0.2                   | 24        | 104            | $7.8\sigma$       |

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| 0.05                  | 23        | 43             | 156        | $2.7\sigma$       |
| 0.1                   | 23        | 63             | 156        | $4.3\sigma$       |
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- Future superbeam experiments using the  $\nu_\mu \rightarrow \nu_e$  channel for hierarchy determination will find it difficult to achieve the required sensitivities due to degeneracies
- It is worthwhile to explore the  $\nu_\mu \rightarrow \nu_\mu$  channel for this purpose since it is largely free of degeneracies.

- Large matter effects are not only confined to  $P_{\mu e}$  but also arise in  $P_{\mu \tau}$  at GeV energies at very long Earth baselines. Both must be properly considered when evaluating the event-rates for experiments measuring muon survival.
- The effects discussed above are significantly sensitive to  $\theta_{13}$  and Sign of  $\Delta m_{31}^2$ , determination of which are outstanding problems of neutrino physics.
- We have tried to show that there is a good possibility that one can determine the Sign of  $\Delta m_{31}^2$  using atmospheric neutrinos in a large mass iron calorimeter with charge identification capabilities.

## Matter effect in $\nu_\mu \rightarrow \nu_\tau$

♠ Click here for  $P_{\mu\tau} \rightsquigarrow$

- ◊ The animation shows the matter effects building up in  $P_{\mu\tau}$  for  $L = 6000$  to  $10500$  km.
- ◊ Maximum matter effects seen at 9700 km.
- ◊ The term wise break up of probability is also depicted in the animation.



To Summary page

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